

Hidden-charm Pentaquark States in Heavy Ion Collisions at the Large Hadron Collider

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In the framework of the quark combination, we derive the yield formulas and study the yield ratios of the hidden-charm pentaquark states in ultra-relativistic heavy ion collisions. We propose some interesting yield ratios which clearly exhibit the production relationships between different hidden-charm pentaquark states. We show how to employ a specific quark combination model to evaluate the yields of exotic $P_c^+(4380)$, $P_c^+(4450)$ and their partners on the basis of reproducing the yields of normal identified hadrons, and execute the calculations in central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV as an example.

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I. INTRODUCTION

The LHCb Collaboration at the Large Hadron Collider (LHC) has recently announced two exotic resonances $P_c^+(4380)$ and $P_c^+(4450)$, consistent with pentaquark states, via the $J/\psi p$ invariant mass spectrum in Λ_b^0 decays in $p+p$ reactions [1]. Thereupon, the study for the internal configurations of such resonances becomes a subject of an intensive discussion in literatures. The so-far suggested interpretations for $P_c^+(4380)$ and $P_c^+(4450)$ include molecular states bound of a charmed baryon and an anticharmed meson [2, 3], molecular states with J/ψ and an excited nucleon [4], composites containing colored baryonlike and mesonlike constituents [5], pentaquark states with one heavy-light diquark, one light-light diquark and a charm antiquark [6–8], pentaquark states with color-antitriplet diquark cu and color-triplet $\bar{c}ud$ [9, 10], and so on. There are also some works about needing further confirmation of whether such resonance is just a kinematical effect or a real exotic resonance by analyzing other processes, such as $\Lambda_b^0 \rightarrow K^- \chi_{c1} p$ [11–14]. Further exploring for the intrinsic dynamical nature of such exotic resonances and searching for their partners are necessary for understanding some longstanding questions in hadronic physics, which requires more theoretical and experimental efforts.

Heavy ion collisions with ultra-relativistic collision energies, especially those at the LHC, provide preferable conditions for the creation of particles in heavy flavor sectors. The studies of heavy flavor exotic resonances in heavy ion collisions not only are complementary to those in the LHCb experiments, but also offer unique insights into some of the fundamental questions in hadronic physics [15, 16]. As is well known, the Quark Combination Mechanism (QCM) is an effective phenomenological method to deal with the hadronization of the partonic system produced early in high energy heavy ion collisions. It has shown successes in reproducing multiplicities, yield ratios, momentum distributions, elliptic flows, etc., of normal identified light, strange and heavy flavor hadrons [17–23], and has also many applications in describing the production properties of exotic resonances [15, 16, 24–27]. Due to the produced bulk deconfined fireball and the hadronization through the quark combination in high energy heavy ion collisions, various kinds of exotic hadrons can be formed. The purpose of this paper is to investigate the yield ratios and estimate the yields of various hidden-charm pentaquark states with different valence quark ingredients and different intrinsic quantum numbers in the framework of the QCM in ultra-relativistic heavy ion collisions at the LHC. This can provide useful references for the exotic hadron search in future experiments and is helpful for understanding the production mechanism of the exotic resonances in heavy ion collisions.

The rest of the paper is organized as follows. In Sec. II, we derive the yield formulas and study the yield ratios of hidden-charm pentaquark states from the basic ideas of the QCM. During the derivation, a few assumptions, approximations and/or simplifications are used and they are all clearly presented. In Sec. III, we show how to employ a specific quark combination model to estimate the yields of different hidden-charm pentaquark states, and give the estimated results in central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV as an example. Sec. IV summaries our work.

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II. YIELD RATIOS OF DIFFERENT HIDDEN-CHARM PENTAQUARK STATES

In this section, we derive the yield formalism of the hidden-charm pentaquark states in the quark combination models based on the basic ideas. We first begin with a quark-antiquark system as general as possible. Then we simplify the results by adopting a few explicit assumptions and/or approximations. Finally we present some interesting results for the yield ratios of different hidden-charm pentaquark states.

A. The yield formalism of the hidden-charm pentaquark states

We start from a color-neutral quark-antiquark system with N_{q_i} quarks of flavor q_i ($q_i = u, d, s, c$) and $N_{\bar{q}_i}$ antiquarks of flavor \bar{q}_i ($\bar{q}_i = \bar{u}, \bar{d}, \bar{s}, \bar{c}$). All these quarks and antiquarks can hadronize via the quark combination into not only normal mesons, baryons and antibaryons, but also exotic hadrons such as tetraquark states, pentaquark states and so on. The momentum distribution $f_{P_j}(p; N_{q_i}, N_{\bar{q}_i})$ for the directly produced pentaquark state P_j with the known quark contents ($q_0 q_1 q_2 q_3 q_4 \bar{q}_5$) after quark combination hadronization is given by

$$f_{P_j}(p; N_{q_i}, N_{\bar{q}_i}) = \sum_{q_1 q_2 q_3 q_4 \bar{q}_5} \int dp_1 dp_2 dp_3 dp_4 dp_5 f_{q_1 q_2 q_3 q_4 \bar{q}_5}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) \mathcal{R}_{P_j, q_1 q_2 q_3 q_4 \bar{q}_5}(p, p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}), \quad (1)$$

where $f_{q_1 q_2 q_3 q_4 \bar{q}_5}$ is the five-particle joint momentum distribution for $(q_1 q_2 q_3 q_4 \bar{q}_5)$. The kernel function $\mathcal{R}_{P_j, q_1 q_2 q_3 q_4 \bar{q}_5}$ stands for the probability density for q_1, q_2, q_3, q_4 , and \bar{q}_5 with momenta p_1, p_2, p_3, p_4 , and p_5 to combine into a pentaquark state P_j of momentum p . Integrating over p from Eq. (1), we can obtain the average number of the directly produced P_j as

$$\bar{N}_{P_j}(N_{q_i}, N_{\bar{q}_i}) = \sum_{q_1 q_2 q_3 q_4 \bar{q}_5} \int dp dp_1 dp_2 dp_3 dp_4 dp_5 f_{q_1 q_2 q_3 q_4 \bar{q}_5}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) \mathcal{R}_{P_j, q_1 q_2 q_3 q_4 \bar{q}_5}(p, p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}). \quad (2)$$

Eqs. (1) and (2) are the most general starting point of describing the production of pentaquark states with any flavors of quark ingredients in high energy reactions based on the basic ideas of the QCM. Different models are some special examples of the general case we consider in these equations. In specific models, different methods and/or assumptions are introduced to construct the precise form of the kernel function. For example, the kernel function evolves into the Wigner function in the coalescence model [20] and the recombination function in the quark recombination model [28], respectively.

For a special kind of hidden-charm pentaquark states $P_j^{c\bar{c}}$ with the known quark contents ($l_0 l_1 l_2 l_3 c \bar{c}$) distinguished by the superscript $c\bar{c}$, we from Eq. (2) easily have

$$\bar{N}_{P_j^{c\bar{c}}}(N_{q_i}, N_{\bar{q}_i}) = \sum_{l_1 l_2 l_3} \int dp dp_1 dp_2 dp_3 dp_4 dp_5 f_{l_1 l_2 l_3 c \bar{c}}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) \mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}(p, p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}), \quad (3)$$

where $l_i = u, d, s$. The joint momentum distribution $f_{l_1 l_2 l_3 c \bar{c}}$ is the number density that satisfies

$$\int dp_1 dp_2 dp_3 dp_4 dp_5 f_{l_1 l_2 l_3 c \bar{c}}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) = N_{l_1 l_2 l_3 c \bar{c}} = N_{l_1 l_2 l_3} N_c N_{\bar{c}}, \quad (4)$$

where $N_{l_1 l_2 l_3 c \bar{c}}$ and

$$N_{l_1 l_2 l_3} = \begin{cases} N_{l_1} N_{l_2} N_{l_3} & \text{for } l_1 \neq l_2 \neq l_3 \\ N_{l_1} (N_{l_1} - 1) N_{l_3} & \text{for } l_1 = l_2 \neq l_3, \\ N_{l_1} (N_{l_1} - 1) (N_{l_1} - 2) & \text{for } l_1 = l_2 = l_3 \end{cases} \quad (5)$$

are the numbers of all the possible $(l_1 l_2 l_3 c \bar{c})$'s and $(l_1 l_2 l_3)$'s, respectively, in the considered bulk quark-antiquark system. We rewrite

$$f_{l_1 l_2 l_3 c \bar{c}}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) = N_{l_1 l_2 l_3 c \bar{c}} f_{l_1 l_2 l_3 c \bar{c}}^{(n)}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}), \quad (6)$$

so that the joint momentum distribution is normalized to one where is denoted by the superscript (n) , i.e.,

$$\int dp_1 dp_2 dp_3 dp_4 dp_5 f_{l_1 l_2 l_3 c \bar{c}}^{(n)}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) = 1. \quad (7)$$

We adopt an assumption of u , d , s -flavor independence of the normalized joint momentum distribution of quarks and/or antiquarks, i.e.,

$$f_{l_1 l_2 l_3 c \bar{c}}^{(n)}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) = f_{llc\bar{c}}^{(n)}(p_1, p_2, p_3, p_4, p_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}), \quad (8)$$

to simplify the formalism. Here, l stands for u , d or s ; N_l and $N_{\bar{l}}$ stand for the total number of u , d and s quarks and that of \bar{u} , \bar{d} and \bar{s} antiquarks, respectively; N_q and $N_{\bar{q}}$ stand for the total number of u , d , s , c quarks and that of \bar{u} , \bar{d} , \bar{s} , \bar{c} antiquarks in the considered quark-antiquark system. With the normalized u , d , s -flavor independent joint momentum distribution, we have

$$\begin{aligned} \bar{N}_{P_j^{c\bar{c}}}(N_{q_i}, N_{\bar{q}_i}) &= \sum_{l_1 l_2 l_3} N_{l_1 l_2 l_3 c \bar{c}} \int dp dp_1 dp_2 dp_3 dp_4 dp_5 f_{llc\bar{c}}^{(n)}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) \\ &\quad \times \mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}(p, p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}). \end{aligned} \quad (9)$$

The kernel function $\mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}$ carries a lot of kinematical and dynamical information on the quark combination hadronization, but its precise form is very ambiguous at present due to its complicated non-perturbative property. Despite all this, we have known the kernel function should be endowed with the following four constraints at least. The first is satisfying the momentum conservation, so it should contain the item $\delta(\sum_{i=1}^5 p_i - p)$. The second is the quark flavor conservation during the hadronization process, which is the requirement of the flavor conservation in the strong interaction. The third is the requirement of the hadronization unitarity, i.e., the production of normal and exotic hadrons should exhaust all quarks and antiquarks in the system after hadronization. The fourth is the dynamics of hadron-hadron production competition, i.e., when a quark hadronizes, whether it forms a meson by combining an antiquark or forms a baryon by combining other two quarks, or goes into an exotic state as other constraints are all satisfied in these cases. Similar detailed discussions of kernel functions on normal hadrons can be found in our previous work [29]. Based on the above four points, we assume a simple but explicit case where the kernel function can be factorized as follows

$$\mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}(p, p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) = C_{P_j^{c\bar{c}}} \mathcal{R}_{l_0 l_1 l_2 l_3} \mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) \delta\left(\sum_{i=1}^5 p_i - p\right), \quad (10)$$

where the δ function guarantees the momentum conservation. $\mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}$ denotes the probability of l_1 , l_2 , l_3 , c , and \bar{c} with momenta p_1 , p_2 , p_3 , p_4 , and p_5 to combine into a hidden-charm pentaquark state $P^{c\bar{c}}$, and it should depend on the momenta of the constituents and their situated environments represented by N_{q_i} and $N_{\bar{q}_i}$. $\mathcal{R}_{l_0 l_1 l_2 l_3}$ guarantees the quark flavor conservation during the hadronization process and the product $\mathcal{R}_{l_0 l_1 l_2 l_3} \mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}$ presents the probability for the l_1 , l_2 , l_3 , c and \bar{c} to combine into a pentaquark state $P^{c\bar{c}}$ with the given quark contents ($l_0 l_1 l_2 l_3 c \bar{c}$). So $\mathcal{R}_{l_0 l_1 l_2 l_3}$ contains the Kronecker δ 's and equals to $N_{iter} \delta_{l_1, l_0} \delta_{l_2, l_0} \delta_{l_3, l_0}$. N_{iter} stands for the number of possible iterations of $l_0 l_1 l_2 l_3$, and it is taken to be 1, 3, and 6 for three identical flavor, two different flavor, and three different flavor cases, respectively. For example, $\mathcal{R}_{uud} = (\delta_{l_1, u} \delta_{l_2, u} \delta_{l_3, d} + \delta_{l_1, u} \delta_{l_2, d} \delta_{l_3, u} + \delta_{l_1, d} \delta_{l_2, u} \delta_{l_3, u})$. $C_{P_j^{c\bar{c}}}$ denotes the probability for that the intrinsic quantum numbers of the formed $P^{c\bar{c}}$ with the quark contents ($l_0 l_1 l_2 l_3 c \bar{c}$) are the same as those of $P_j^{c\bar{c}}$. A pentaquark state, composed of five quarks and antiquarks, every of which has spin 1/2, has three different spin quantum numbers, i.e., $J = 1/2, 3/2, 5/2$. We introduce $R_{J_{31}}$ and $R_{J_{51}}$ to denote the relative production ratio of $J = 3/2$ to $J = 1/2$ pentaquark states and that of $J = 5/2$ to $J = 1/2$ pentaquark states with the same flavor compositions. Considering that each quark has parity +1 and each antiquark has parity -1, and we only include the ground $L = 0$ and the first excited $L = 1$ pentaquark states, a pentaquark state has two different parity quantum numbers, i.e., $P = -1, +1$. We introduce $R_{P_{10}}$ to denote the relative production ratio of $L = 1$ to $L = 0$ pentaquark states with the same flavor compositions. So we can obtain

$$C_{P_j^{c\bar{c}}} = \begin{cases} \frac{1}{(1+R_{J_{31}}+R_{J_{51}})} \times \frac{1}{(1+R_{P_{10}})} & \text{for } J^P = (1/2)^- \text{ states} \\ \frac{1}{(1+R_{J_{31}}+R_{J_{51}})} \times \frac{R_{P_{10}}}{(1+R_{P_{10}})} & \text{for } J^P = (1/2)^+ \text{ states} \\ \frac{R_{J_{31}}}{(1+R_{J_{31}}+R_{J_{51}})} \times \frac{1}{(1+R_{P_{10}})} & \text{for } J^P = (3/2)^- \text{ states} \\ \frac{R_{J_{31}}}{(1+R_{J_{31}}+R_{J_{51}})} \times \frac{R_{P_{10}}}{(1+R_{P_{10}})} & \text{for } J^P = (3/2)^+ \text{ states} \\ \frac{R_{J_{51}}}{(1+R_{J_{31}}+R_{J_{51}})} \times \frac{1}{(1+R_{P_{10}})} & \text{for } J^P = (5/2)^- \text{ states} \\ \frac{R_{J_{51}}}{(1+R_{J_{31}}+R_{J_{51}})} \times \frac{R_{P_{10}}}{(1+R_{P_{10}})} & \text{for } J^P = (5/2)^+ \text{ states.} \end{cases} \quad (11)$$

$R_{J_{31}}$ and $R_{J_{51}}$ have been determined to be 1.6 and 0.6 according to a simple spin counting. $R_{P_{10}}$ has been determined to be 0.258 by the Wigner function method in Ref. [15]. The still left unknown item is $\mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i})$. Recalling that it is the probability of l_1 , l_2 , l_3 , c , and \bar{c} with momenta p_1 , p_2 , p_3 , p_4 , and p_5 to combine into a hidden-charm pentaquark state $P^{c\bar{c}}$, we assume that it is u , d , s -flavor independent, i.e.,

$$\mathcal{R}_{P_j^{c\bar{c}}, l_1 l_2 l_3 c \bar{c}}(p_1, p_2, p_3, p_4, p_5; N_{q_i}, N_{\bar{q}_i}) = \mathcal{R}_{P_j^{c\bar{c}}, llc\bar{c}}(p_1, p_2, p_3, p_4, p_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}). \quad (12)$$

Substitute Eqs. (10) and (12) into Eq. (9), we have

$$\begin{aligned} \bar{N}_{P_j^{c\bar{c}}}(N_{q_i}, N_{\bar{q}_i}) &= \sum_{l_1 l_2 l_3} N_{l_1 l_2 l_3 c\bar{c}} C_{P_j^{c\bar{c}}} \mathcal{R}_{l_{01} l_{02} l_{03}} \\ &\times \int dp dp_1 dp_2 dp_3 dp_4 dp_5 f_{llc\bar{c}}^{(n)}(p_1, p_2, p_3, p_4, p_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) \mathcal{R}_{P_j^{c\bar{c}}, llc\bar{c}}(p_1, p_2, p_3, p_4, p_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) \delta\left(\sum_{i=1}^5 p_i - p\right). \end{aligned} \quad (13)$$

We denote the momentum integral in the second line in Eq. (13) to be $\gamma_{P_j^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}})$ and then obtain

$$\bar{N}_{P_j^{c\bar{c}}}(N_{q_i}, N_{\bar{q}_i}) = \sum_{l_1 l_2 l_3} N_{l_1 l_2 l_3 c\bar{c}} C_{P_j^{c\bar{c}}} \mathcal{R}_{l_{01} l_{02} l_{03}} \gamma_{P_j^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}}). \quad (14)$$

Assuming that the internal configurations of all different species of hidden-charm pentaquark states are the same and summing over different species of hidden-charm pentaquark states, we obtain the average number of all the hidden-charm pentaquark states $\bar{N}_{P^{c\bar{c}}}$ as follows

$$\bar{N}_{P^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) = N_{llc\bar{c}} \gamma_{P^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}}), \quad (15)$$

where $N_{llc\bar{c}} = N_{lll} N_c N_{\bar{c}} = N_l(N_l - 1)(N_l - 2)N_c N_{\bar{c}}$. Substitute Eq. (15) into Eq. (14), the average number of a specified hidden-charm pentaquark state $P_j^{c\bar{c}}$ is given by

$$\begin{aligned} \bar{N}_{P_j^{c\bar{c}}}(N_{q_i}, N_{\bar{q}_i}) &= \sum_{l_1 l_2 l_3} C_{P_j^{c\bar{c}}} \mathcal{R}_{l_{01} l_{02} l_{03}} \frac{N_{l_1 l_2 l_3 c\bar{c}}}{N_{llc\bar{c}}} \bar{N}_{P^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) \\ &= C_{P_j^{c\bar{c}}} N_{iter} \frac{N_{l_{01} l_{02} l_{03} c\bar{c}}}{N_{llc\bar{c}}} \bar{N}_{P^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) \\ &= C_{P_j^{c\bar{c}}} N_{iter} \frac{N_{l_{01} l_{02} l_{03}}}{N_{lll}} \bar{N}_{P^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}}). \end{aligned} \quad (16)$$

For a reaction at a given energy, the average numbers of quarks of different flavors $\langle N_{q_i} \rangle$ and those of antiquarks of different flavors $\langle N_{\bar{q}_i} \rangle$ are fixed while N_{q_i} and $N_{\bar{q}_i}$ follow a certain fluctuation distribution. In this work, we focus on the midrapidity region at high LHC energy where the influence of net quarks from the colliding nuclei is negligible [30]. We suppose a polynomial distribution for both the numbers of u , d and s quarks at a given N_l and the numbers of \bar{u} , \bar{d} and \bar{s} at a given $N_{\bar{l}}$ with the prior probabilities $p_u = p_d = p_{\bar{u}} = p_{\bar{d}} = 1/(2 + \lambda_s)$, $p_s = p_{\bar{s}} = \lambda_s/(2 + \lambda_s)$. Here, we introduce λ_s to denote the production suppression of strange quarks. Averaging over this distribution, Eq. (16) becomes

$$\bar{N}_{P_j^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) = C_{P_j^{c\bar{c}}} N_{iter} p_{l_{01}} p_{l_{02}} p_{l_{03}} \bar{N}_{P^{c\bar{c}}}(N_l, N_{\bar{l}}, N_q, N_{\bar{q}}), \quad (17)$$

We will also consider the fluctuations of N_l , $N_{\bar{l}}$, N_q and $N_{\bar{q}}$ in the given kinematic region. By averaging over this fluctuation distribution with the fixed $\langle N_l \rangle$, $\langle N_{\bar{l}} \rangle$, $\langle N_q \rangle$ and $\langle N_{\bar{q}} \rangle$, we have

$$\langle N_{P_j^{c\bar{c}}} \rangle (\langle N_l \rangle, \langle N_{\bar{l}} \rangle, \langle N_q \rangle, \langle N_{\bar{q}} \rangle) = C_{P_j^{c\bar{c}}} N_{iter} p_{l_{01}} p_{l_{02}} p_{l_{03}} \langle N_{P^{c\bar{c}}} \rangle (\langle N_l \rangle, \langle N_{\bar{l}} \rangle, \langle N_q \rangle, \langle N_{\bar{q}} \rangle), \quad (18)$$

where $\langle N_{P^{c\bar{c}}} \rangle$ stands for the total average number of all the hidden-charm pentaquark states produced in the combination process. Eq. (18) is the general yield formulism of hidden-charm pentaquark states obtained from the basic ideas of the QCM with three assumptions, i.e., the u , d , s flavor independence of the normalized quark/antiquark joint momentum distribution, the factorization of the kernel function and the u , d , s flavor independence of the momentum dependent part $\mathcal{R}_{P_j^{c\bar{c}}, llc\bar{c}}$, and the same configuration for all different species of hidden-charm pentaquark states.

B. Yield ratios of different hidden-charm pentaquark states

From the yield formula in Eq. (18), we see that there exist many simple relationships between the yields of different hidden-charm pentaquark states. It is necessary to stress that these relations are the general features of the QCM under the three assumptions mentioned in the last subsection. Actrally, they are independent of the detailed form of the momentum dependence of the kernel function or the momentum distributions of the quarks and antiquarks. They even do not depend on whether other exotic states, such as tetraquark states, pentaquark states in light sectors and so on, are produced during the hadronization process. They are the characteristics for hidden-charm pentaquark resonance production in the QCM, which can be used to test the combination production mechanism of the exotic pentaquark states. We will propose some of these interesting relations in

the following. The first group concerns hidden-charm pentaquark states with the same spin and parity quantum numbers, and they are listed as follows

$$\frac{\langle N_{P^{c\bar{c}}(uudc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uddc\bar{c})} \rangle} = 1, \quad \frac{\langle N_{P^{c\bar{c}}(uuuc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(dddc\bar{c})} \rangle} = 1, \quad (19)$$

$$\frac{\langle N_{P^{c\bar{c}}(uusc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(ddsc\bar{c})} \rangle} = 1, \quad \frac{\langle N_{P^{c\bar{c}}(ussc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(dssc\bar{c})} \rangle} = 1, \quad (20)$$

$$\frac{\langle N_{P^{c\bar{c}}(udsc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uuuc\bar{c})} \rangle} = 2, \quad \frac{\langle N_{P^{c\bar{c}}(uudc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uuuc\bar{c})} \rangle} = 3. \quad (21)$$

In the above the quark contents in the parentheses denote the flavor ingredients of the corresponding pentaquark states. These ratios are independent of the collision energy, the species of the colliding nuclei and the collision centrality. Therefore, they can be used to test the universality of the production mechanism of four quarks and an antiquark combining into a pentaquark state. In addition, they clearly exhibit relative production weights for different pentaquark states, which can be used by the experimental search for more pentaquark states. For example, if the resonances observed by the LHCb Collaboration are indeed pentaquark states with $(uudc\bar{c})$ quark contents and the hidden-charm pentaquark states are indeed produced via the quark combination, we have no reason to say there are no $P^{c\bar{c}}(uddc\bar{c})$ produced even it has not been observed simultaneously.

The second group is related with the strangeness production denoted by λ_s , and they are as

$$\frac{\langle N_{P^{c\bar{c}}(uusc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uudc\bar{c})} \rangle} = \lambda_s, \quad \frac{\langle N_{P^{c\bar{c}}(ussc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uudc\bar{c})} \rangle} = \lambda_s^2, \quad (22)$$

$$\frac{\langle N_{P^{c\bar{c}}(ussc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uusc\bar{c})} \rangle} = \lambda_s, \quad \frac{\langle N_{P^{c\bar{c}}(sssc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uusc\bar{c})} \rangle} = \frac{\lambda_s^2}{3}, \quad (23)$$

$$\frac{\langle N_{P^{c\bar{c}}(sssc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(ussc\bar{c})} \rangle} = \frac{\lambda_s}{3}, \quad \frac{\langle N_{P^{c\bar{c}}(sssc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(udsc\bar{c})} \rangle} = \frac{\lambda_s^2}{6}, \quad (24)$$

$$\frac{\langle N_{P^{c\bar{c}}(udsc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uudc\bar{c})} \rangle} = 2\lambda_s, \quad \frac{\langle N_{P^{c\bar{c}}(udsc\bar{c})} \rangle}{\langle N_{P^{c\bar{c}}(uuuc\bar{c})} \rangle} = 6\lambda_s. \quad (25)$$

These ratios are also the general results of the quark combination with the above three assumptions, and they show the relative production weights of strange hidden-charm pentaquark states to those without strangeness. Generally speaking, the production of strange hidden-charm pentaquark states is suppressed relative to non-strange ones due to the strangeness production suppression. This just shows the advantages of heavy ion collisions to produce strange hidden-charm pentaquark states compared to elementary particle reactions because of the strangeness enhancement in high energy heavy ion collisions. More interestingly, from Eq. (25) one can see that strange $P^{c\bar{c}}(udsc\bar{c})$ is nearly not suppressed relative to $P^{c\bar{c}}(uudc\bar{c})$ and instead enhanced by a factor of about three relative to the non-strange $P^{c\bar{c}}(uuuc\bar{c})$ considering that λ_s in heavy ion collisions is located in the range (0.4-0.5) [31]. These results provide important references for future experimental search for strange hidden-charm pentaquark states.

As a brief summary of Sec. II, we want to emphasize once more that the method we consider in this section is intended to be a general case based on the basic ideas of the QCM. To simplify the results, we adopt a few assumptions and/or approximations based on symmetry and general principles, such as the u , d , s flavor independence of the normalized quark/antiquark joint momentum distribution, the factorization of the kernel function and the u , d , s flavor independence of the momentum dependent part $\mathcal{R}_{P^{c\bar{c}},lllc\bar{c}}$, and the same configuration for all different species of hidden-charm pentaquark states. Except the assumptions and/or approximations stated explicitly in the work, the results do not depend on other factors. They do not depend on the specific collision energy, colliding nuclei or the centrality. They even are not limited to heavy ion collisions, but can be suitable in elementary particle reactions if the reaction energy is large enough to create a relative bulk quark-antiquark system, e.g., possibly in LHC $p + p$ reactions. Therefore, these results on yield ratios can be used to test the combination production mechanism of the pentaquark states. We focus on different hidden-charm pentaquark state yield ratios in this section. No efforts are made to study the yields, and we leave them to the next section.

III. ESTIMATES FOR THE YIELDS OF HIDDEN-CHARM PENTAQUARK STATES

The yields of various exotic resonances are the most fundamental quantities and valuable probes for exploring their internal configurations and their production mechanisms [15, 16, 26, 27, 32]. Phenomenological models such as the statistical model and the coalescence model have been used to predict the yields for different possible exotic states, e.g., $f_0(980)$, $a_0(980)$, $X(3872)$, $D_{sJ}(2317)$, $\Theta_{cs}(uudsc\bar{c})$, etc., [15, 16, 26, 27]. In this section, we will evaluate the yields for different species of the hidden-charm pentaquark resonances produced in heavy ion collisions at the LHC.

Recalling Eq. (13), the average number of the produced $P_j^{c\bar{c}}$ can be written after integrating over p as follows

$$\begin{aligned} \bar{N}_{P_j^{c\bar{c}}}(N_{q_i}, N_{\bar{q}_i}) &= \sum_{l_1 l_2 l_3} N_{l_1 l_2 l_3 c\bar{c}} C_{P_j^{c\bar{c}}} \mathcal{R}_{l_{01} l_{02} l_{03}} \\ &\times \int dp_1 dp_2 dp_3 dp_4 dp_5 f_{llc\bar{c}}^{(n)}(p_1, p_2, p_3, p_4, p_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) \mathcal{R}_{P_j^{c\bar{c}}, llc\bar{c}}(p_1, p_2, p_3, p_4, p_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}). \end{aligned} \quad (26)$$

To further simplify the formalism, we ignore the transverse momentum information and only include the longitudinal part. We adopt the rapidity coordinate y to replace the momentum coordinate and obtain

$$\begin{aligned} \bar{N}_{P_j^{c\bar{c}}}(N_{q_i}, N_{\bar{q}_i}) &= \sum_{l_1 l_2 l_3} N_{l_1 l_2 l_3 c\bar{c}} C_{P_j^{c\bar{c}}} \mathcal{R}_{l_{01} l_{02} l_{03}} \\ &\times \int dy_1 dy_2 dy_3 dy_4 dy_5 f_{llc\bar{c}}^{(n)}(y_1, y_2, y_3, y_4, y_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) \mathcal{R}_{P_j^{c\bar{c}}, llc\bar{c}}(y_1, y_2, y_3, y_4, y_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}). \end{aligned} \quad (27)$$

Confining ourselves to the hadron production at the midrapidity $y \in (-0.5, 0.5)$, we set a uniform quark joint distribution, i.e., $f_{llc\bar{c}}^{(n)}(y_1, y_2, y_3, y_4, y_5; N_l, N_{\bar{l}}, N_q, N_{\bar{q}}) = 1$. The item $\mathcal{R}_{P_j^{c\bar{c}}, llc\bar{c}}$ has no analytical form currently, so we rely on a specific Quark Combination Model developed by ShanDong group SDQCM [33, 34].

In the following, we first give a brief introduction to the old version of the SDQCM [24, 33, 34] and show how to extend the old version to include the production of the exotic multiquark states. Then we use the SDQCM (old version and new version) to calculate the yields of the normal identified light, strange and charm hadrons. Finally, we estimate the yields of hidden-charm pentaquark states. The results in central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV are given as an example.

A. An introduction to the SDQCM

The starting point of SDQCM is a color singlet system that consists of constituent quarks and antiquarks. All of these quarks and antiquarks are lined up in phase space, e.g., in a one-dimensional rapidity axis, and then combine into different initial hadrons one by one according to the following quark combination rule.

- (i) Start from the first parton (q_1 or \bar{q}_1) in the line.
- (ii) If the baryon number of the second parton in the line is different from that of the first, i.e., the first two partons are either $q_1 \bar{q}_2$ or $\bar{q}_1 q_2$, they combine into a meson and are removed from the line; then go back to point (i). If they are either $q_1 q_2$ or $\bar{q}_1 \bar{q}_2$, then go to the next point.
- (iii) Look at the third parton, if its baryon number is the same as the first two (i.e., $q_1 q_2 q_3$ or $\bar{q}_1 \bar{q}_2 \bar{q}_3$), the first three partons combine into a baryon or an antibaryon and are removed from the line; then go back to point (i). Otherwise, (i.e., $q_1 q_2 \bar{q}_3$ or $\bar{q}_1 \bar{q}_2 q_3$), considering the color factor for quark-antiquark pairs in the color singlet channel is about twice larger than that for diquarks in the color antitriplet, $q_1 \bar{q}_3$ or $\bar{q}_1 q_3$ form a meson and are removed from the line; q_2 or \bar{q}_2 is still left as the first parton in the line; then go back to point (i).

Here gives an example to show how the above quark combination rule works

$$\begin{aligned} &q_1 \bar{q}_2 \bar{q}_3 \bar{q}_4 \bar{q}_5 \bar{q}_6 \bar{q}_7 q_8 q_9 \bar{q}_{10} q_{11} q_{12} q_{13} q_{14} q_{15} q_{16} \bar{q}_{17} q_{18} q_{19} \cdots \cdots \\ &\rightarrow M(q_1 \bar{q}_2) \bar{B}(\bar{q}_3 \bar{q}_4 \bar{q}_5) M(\bar{q}_6 q_8) M(\bar{q}_7 q_9) M(\bar{q}_{10} q_{11}) B(q_{12} q_{13} q_{14}) M(q_{15} \bar{q}_{17}) B(q_{16} q_{18} q_{19}) \cdots \cdots \end{aligned} \quad (28)$$

After applying the quark combination rule in the quark-antiquark system, the numbers of mesons, baryons and antibaryons with given quark contents are determined. Including the production weights of different hadrons with the same quark contents which are determined by their intrinsic quantum numbers such as spin and parity, the multiplicities of various kinds of initial hadrons are obtained. A more detailed description for the old version of the SDQCM can be found in Refs. [24, 33, 34].

The above old version of SDQCM sets the priority of the smallest number of partons to form a hadron, and there are no multiquark states produced. In the following, we will extend the quark combination rule to include exotic multiquark states. Considering the large difference between heavy charm quarks and u , d , s quarks, the previous quark combination rule is extended to be a new one as follows.

- (i) Start from the first parton (q_1 or \bar{q}_1) in the line.
- (ii) If the baryon number of the second parton in the line is different from that of the first, i.e., the first two partons are either $q_1 \bar{q}_2$ or $\bar{q}_1 q_2$, they combine into a meson and are removed from the line; then go back to point (i). If they are either $q_1 q_2$ or $\bar{q}_1 \bar{q}_2$, then go to point (iii).

- (iii) Look at the third parton, if its baryon number is the same as the first two, (i.e., $q_1q_2q_3$ or $\bar{q}_1\bar{q}_2\bar{q}_3$), the first three partons combine into a baryon or an antibaryon and are removed from the line; then go back to point (i). Otherwise, (i.e., $q_1q_2\bar{q}_3$ or $\bar{q}_1\bar{q}_2q_3$), if the first and the third partons are both heavy charm flavors or light/strange flavors, $q_1\bar{q}_3$ or \bar{q}_1q_3 form a meson and are removed from the line; q_2 or \bar{q}_2 is still left in the line; then go back to point (i). If one of the first and the third partons is in charm flavor and the other is in light/strange flavors, the probability for them to combine into a meson will decrease due to their relatively large momentum difference resulted from their large mass difference [20]. If they can not combine into D mesons, multiquark states are the most possible candidates for them to contribute to. We use ε to denote the probability for the first and the third partons to combine with other nearby partons to merge into multiquark states, and in this case, go to point (iv).
- (iv) Look at the fourth parton, if its baryon number is the same as the third, (i.e., $q_1q_2\bar{q}_3\bar{q}_4$ or $\bar{q}_1\bar{q}_2q_3q_4$), the first four partons combine into a tetraquark state, and then go back to point (i). Otherwise, (i.e., $q_1q_2\bar{q}_3q_4$ or $\bar{q}_1\bar{q}_2q_3\bar{q}_4$), go to point (v).
- (v) Look at the fifth parton, if its baryon number is the same as the fourth, (i.e., $q_1q_2\bar{q}_3q_4q_5$ or $\bar{q}_1\bar{q}_2q_3\bar{q}_4\bar{q}_5$), the first five partons combine into a pentaquark state, and then go back to point (i). Otherwise, (i.e., $q_1q_2\bar{q}_3q_4\bar{q}_5$ or $\bar{q}_1\bar{q}_2q_3\bar{q}_4q_5$), go to point (vi).
- (vi) Look at the sixth parton, if its baryon number is the same as the fifth, (i.e., $q_1q_2\bar{q}_3q_4\bar{q}_5\bar{q}_6$ or $\bar{q}_1\bar{q}_2q_3\bar{q}_4q_5q_6$), the first six partons combine into a sixquark state, and then go back to point (i). Otherwise, continue to look at the next parton or the next next parton until they can be in a color-singlet multiquark state.

In the following, we give an example to show how the above new quark combination rule works

$$\begin{aligned}
 & q_1\bar{q}_2\bar{q}_3\bar{q}_4\bar{q}_5\bar{l}_6\bar{l}_7l_8l_9\bar{q}_{10}q_{11}q_{12}q_{13}q_{14}l_{15}q_{16}\bar{c}q_{18}q_{19} \cdots \cdots \\
 \rightarrow & \begin{cases} M(q_1\bar{q}_2)\bar{B}(\bar{q}_3\bar{q}_4\bar{q}_5)M(\bar{l}_6l_8)M(\bar{l}_7l_9)M(\bar{q}_{10}q_{11})B(q_{12}q_{13}q_{14})M(l_{15}\bar{c})B(q_{16}q_{18}q_{19}) \cdots \cdots & \text{with probability } 1 - \varepsilon \\ M(q_1\bar{q}_2)\bar{B}(\bar{q}_3\bar{q}_4\bar{q}_5)M(\bar{l}_6l_8)M(\bar{l}_7l_9)M(\bar{q}_{10}q_{11})B(q_{12}q_{13}q_{14})P(l_{15}q_{16}\bar{c}q_{18}q_{19}) \cdots \cdots & \text{with probability } \varepsilon. \end{cases} \quad (29)
 \end{aligned}$$

Note that when $\varepsilon = 0$, the new quark combination rule becomes to be the old one, and when $\varepsilon = 1$, the multiquark states are allowed to be produced most abundantly. Applying the new quark combination rule in the quark-antiquark system, the numbers of mesons, baryons, antibaryons and multiquark states with given quark contents are determined. Including the production weights of different hadrons, the multiplicities of various kinds of initial hadrons are obtained.

We want to state that the new quark combination rule includes the production of the multiquark states such as the tetraquark states, the pentaquark states, the sixquark states and so on in the charm sector. The production of the multiquark states in light and strange sectors are simply neglected because the $q_1q_2\bar{q}_3$ or $\bar{q}_1\bar{q}_2q_3$ with light and/or strange flavors in the line are easily to form mesons. In addition, the production of the multiquark states with more valence quarks and/or antiquarks is suppressed more strongly in the new quark combination rule, which follows the general principles during the hadron production. The quark combination rule in the SDQCM naturally satisfy the near phase space correlation for quarks and/or antiquarks when they combine into different hadrons and the unitary requirement simultaneously. Also, it determines the competition of the hadron production uniquely. Although very simple for the quark combination rule, it can meet the constraints for the kernel functions of hadrons and make the calculations possible.

B. Yields of normal identified hadrons

In this subsection, we use the SDQCM to compute the midrapidity yields dN/dy of normal identified mesons and baryons in light, strange and charm sectors in central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. Before doing that, we have to determine the input parameters of the SDQCM. The first is the number of all the midrapidity quarks and antiquarks $dN_{q+\bar{q}}/dy$ which determines the total multiplicity at midrapidity and is related with the reaction energy, the species of the collision nuclei, the collision centrality, the studied dynamical space, and so on. In this paper, we set $dN_{q+\bar{q}}/dy = 2dN_q/dy = 2dN_{\bar{q}}/dy = 3300$, where the net quark number is set to be zero due to focusing on the midrapidity region at so high LHC energy [30]. The second input parameter is the strangeness suppression factor λ_s that determines the number of the strange quarks. We adopt the saturated value for λ_s in relativistic heavy ion collisions, i.e., $\lambda_s = 0.41$ [29]. The last input parameter is the number of charm quarks in the quark-antiquark system. We set it by extrapolating $p + p$ reaction data at LHC as follows

$$\frac{dN_c}{dy} = \langle T_{AA} \rangle \frac{d\sigma_c^{pp}}{dy} = \langle T_{AA} \rangle \frac{1}{R} \frac{d\sigma_{D^0}^{pp}}{dy} = 21. \quad (30)$$

Here $R = 0.54 \pm 0.05$ is the branch ratio of charm quarks into final D^0 mesons measured in e^+e^- reactions [35]. $\langle T_{AA} \rangle = 26.4 \pm 0.5 \text{ mb}^{-1}$ is the average nuclear overlap function in the most central (0-5% centrality) Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV calculated with the Glauber model [36]. The differential cross section of D^0 is $d\sigma_{D^0}^{pp}/dy = 0.428 \pm 0.115 \text{ mb}$ in $p + p$ reactions at $\sqrt{s} = 2.76 \text{ TeV}$ [37].

TABLE I: Midrapidity yields dN/dy of normal identified mesons and baryons in central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. The experimental data are from Refs. [30, 38–40].

Hadron	Data	SDQCM ($\varepsilon = 0$)	SDQCM ($\varepsilon = 1$)	SDQCM ($\varepsilon = 0.6$)
π^+	733 ± 54	742	728	734
π^-	732 ± 52	742	728	734
K^+	109 ± 9	115	111	112
K^-	109 ± 9	115	111	112
K_S^0	110 ± 10	109	106	107
ϕ	$13.8 \pm 0.5 \pm 1.7$	14.8	14.5	14.6
p	34 ± 3	32	31	32
\bar{p}	32 ± 3	32	31	32
Λ	26 ± 3	25	25	25
$\bar{\Lambda}$	—	25	25	25
Ξ^-	$3.34 \pm 0.06 \pm 0.24$	3.89	3.84	3.86
$\bar{\Xi}^+$	$3.28 \pm 0.06 \pm 0.23$	3.89	3.84	3.86
Ω^-	$0.58 \pm 0.04 \pm 0.09$	0.55	0.54	0.54
$\bar{\Omega}^+$	$0.60 \pm 0.05 \pm 0.09$	0.55	0.54	0.54
$D^+ (D^-)$	—	3.35	2.29	2.71
$D^0 (\bar{D}^0)$	—	10.4	7.10	8.41
$D^{*+} (D^{*-})$	—	5.15	3.52	4.17
$D^{*0} (\bar{D}^{*0})$	—	5.15	3.52	4.17
$D_s^+ (D_s^-)$	—	2.82	1.93	2.28
$D_s^{*+} (D_s^{*-})$	—	2.11	1.44	1.71
η_c	—	0.054	0.055	0.054
J/Ψ	—	0.160	0.166	0.164
$\Lambda_c^+ (\bar{\Lambda}_c^-)$	—	2.86	2.83	2.84
$\Sigma_c^+ (\bar{\Sigma}_c^-)$	—	0.569	0.564	0.566
$\Sigma_c^0 (\bar{\Sigma}_c^0)$	—	0.473	0.469	0.471
$\Sigma_c^{++} (\bar{\Sigma}_c^{--})$	—	0.474	0.470	0.471
$\Xi_{cc}^{++} (\bar{\Xi}_{cc}^{--})$	—	0.0140	0.0139	0.0140
$\Xi_{cc}^+ (\bar{\Xi}_{cc}^-)$	—	0.0140	0.0139	0.0140
$\Omega_{cc}^+ (\bar{\Omega}_{cc}^-)$	—	0.00580	0.00577	0.00574
$\Omega_{ccc}^{++} (\bar{\Omega}_{ccc}^{--})$	—	0.000195	0.000192	0.000193

The calculated midrapidity yields dN/dy of different identified hadrons are collected in Table I. The contributions from strong and electromagnetic decays for light and strange hadrons have been included to coincide with the ALICE experiments. For charmed hadrons, we only consider the decays of D^* mesons for D mesons and Σ_c and Σ_c^* baryons for Λ_c^+ . The third and the fourth columns in Table I are the results calculated by the old version of the SDQCM in which there are no multiquark states included and those by the new version of the SDQCM with $\varepsilon = 1$ where the multiquark states are allowed to be produced most abundantly. From Table I, one can see that the results for light and strange hadrons in these two limit cases, i.e., $\varepsilon = 0$ and $\varepsilon = 1$, are nearly the same and they both agree well with the available data from Refs. [30, 38–40]. For charmed baryons and hidden-charm mesons, the results with $\varepsilon = 1$ and $\varepsilon = 0$ are comparable. For open charmed mesons, the results with $\varepsilon = 1$ decrease about 30% compared to those with $\varepsilon = 0$. The predicted yield of Ω_{ccc}^{++} is in good agreement with the result from the coalescence model [41]. The computed results for the other charmed hadrons wait for the comparisons with the future experimental measurements and/or other theoretical calculations.

C. Estimates for the yields of various hidden-charm pentaquark states

Based on the success in the description of the yields of normal light and strange hadrons, we will in this subsection employ SDQCM to give an estimate for the midrapidity yields of different hidden-charm pentaquark states in the most central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. Setting $\varepsilon = 1$, we obtain the up limit value for the number of all the hidden-charm pentaquark states $dN_{pc\bar{c}}/dy = 0.0411$. Taking it into Eq. (18), we can predict the up limit values for the multiplicities of different hidden-charm pentaquark states. The results are in Table II.

TABLE II: Up limit values for the midrapidity yields dN/dy of various hidden-charm pentaquark states with different spins and parities in the most central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV.

Flavor	$uudc\bar{c}$						$uddc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	21.9	5.65	35.0	9.03	13.1	3.39	21.9	5.65	35.0	9.03	13.1	3.39
Flavor	$uuuc\bar{c}$						$dddc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	7.29	1.88	11.7	3.01	4.38	1.13	7.29	1.88	11.7	3.01	4.38	1.13
Flavor	$uusc\bar{c}$						$ddsc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	8.97	2.31	14.4	3.70	5.38	1.39	8.97	2.31	14.4	3.70	5.38	1.39
Flavor	$ussc\bar{c}$						$dssc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	3.68	0.949	5.89	1.52	2.21	0.569	3.68	0.949	5.89	1.52	2.21	0.569
Flavor	$udsc\bar{c}$						$sssc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	17.9	4.63	28.7	7.41	10.8	2.78	0.503	0.130	0.804	0.208	0.302	0.0778

To further evaluate the yields of different hidden-charm pentaquark states, we have to determine ε . Here we use the experimental data on the p_T distribution of D^0 mesons in the 0-10% centrality measured by the ALICE Collaboration [42]. We use the Blast Wave Model [43] to extract the midrapidity yield for D^0 mesons as about $dN/dy = 7.42$. Simply scaled by the $\langle T_{AA} \rangle = 26.4 \pm 0.5 \text{ mb}^{-1}$ [36] and $\langle T_{AA} \rangle = 23.44 \pm 0.76 \text{ mb}^{-1}$ [42] in the 0-5% centrality and 0-10% centrality, respectively, we have the experimental result $dN/dy \approx 8.36$ for D^0 mesons in the most central 0-5% collisions. By reproducing D^0 mesons, ε is set to be 0.6. The midrapidity yields for other normal identified hadrons when $\varepsilon = 0.6$ are also listed in Table I. With $\varepsilon = 0.6$, we obtain the estimated number of all the hidden-charm pentaquark states is $dN_{pc\bar{c}}/dy = 0.0248$. Taking it into Eq. (18), we can predict the multiplicities for different hidden-charm pentaquark states. The results are in Table III.

Note that we neglect the contribution from the hadron-hadron rescattering in hadronic phase and assume the dominant source is from the quark combination. The pentaquark states made by the quark contents ($uudc\bar{c}$) with $J = 3/2$ and $J = 5/2$ in Table II and Table III correspond to the observed $P_c^+(4380)$ and $P_c^+(4450)$, respectively, from which one can see that the state with

TABLE III: Estimates for the midrapidity yields dN/dy of various pentaquark states with different spins and parities in the most central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV.

Flavor	$uudc\bar{c}$						$uddc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	13.2	3.41	21.1	5.45	7.92	2.04	13.2	3.41	21.1	5.45	7.92	2.04
Flavor	$uuuc\bar{c}$						$dddc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	4.40	1.14	7.04	1.82	2.64	0.681	4.40	1.14	7.04	1.82	2.64	0.681
Flavor	$uusc\bar{c}$						$ddsc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	5.41	1.40	8.66	2.23	3.25	0.838	5.41	1.40	8.66	2.23	3.25	0.838
Flavor	$ussc\bar{c}$						$dssc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	2.22	0.573	3.55	0.916	1.33	0.344	2.22	0.573	3.55	0.916	1.33	0.344
Flavor	$udsc\bar{c}$						$sssc\bar{c}$					
J^P	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺	(1/2) ⁻	(1/2) ⁺	(3/2) ⁻	(3/2) ⁺	(5/2) ⁻	(5/2) ⁺
$\frac{dN}{dy} (10^{-4})$	10.8	2.79	17.3	4.47	6.50	1.68	0.303	0.0783	0.485	0.125	0.182	0.0470

negative parity has a larger yield than that with positive parity by a factor of about four, indicating that the yield measurements can help to determine the parities of such observed resonances. Our results show that the estimated yields of these different hidden-charm pentaquark states are large enough for carrying out realistic measurements.

IV. SUMMARY

With the basic ideas of the QCM and a few assumptions and/or simplifications based on symmetry and general principles, we have derived the yield formulas and studied the yield ratios of different species of hidden-charm pentaquark states in heavy ion collisions at ultra-relativistic collision energies. We found some interesting relations between different hidden-charm pentaquark states. These results are properties of the pentaquark state production in the QCM under these assumptions such as the u , d , s flavor independence of the normalized joint momentum distributions, the factorization of the kernel function and the u , d , s flavor independence of the momentum dependent part of the kernel function, and the same internal configuration for all different species of hidden-charm pentaquark states. They are independent of the particular quark combination models, so they can be used to test the quark combination production mechanism of the exotic resonances. We have also calculated the yields of normal identified light, strange and charm hadrons with the SDQCM, and on this basis we have shown how to estimate the values for the yields of different hidden-charm pentaquark states. The estimated results in the most central Pb+Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV are presented as an example. All of this can help to search for the partners of the observed $P_c^+(4380)$ and $P_c^+(4450)$ and shed light on the understanding of the production mechanism of the exotic pentaquark states in heavy ion collisions.

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